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SUPERNOVA NEUTRINOS: WHAT CAN WE LEARN?

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SUPERNOVA NEUTRINOS



Production (flavor)

- $\langle \psi_i$
- Core-collapse simulation
- Microphysics of SN core
- Stellar nucleosynthesis
- Exotic particles emission

Propagation (mass,mixing) $\int \exp(-iHt)$

- Flavor conversions
- Matter effects: shock wave,turbulences
- Dense neutrino bkg
- New interactions
- Decays

Detection (flavor) $\psi_{f}
angle$

- Interaction cross sections
- Different detection strategies
- Observable signatures

SN 1054 \rightarrow Crab Nebula

Supernova 1054 Petrograph



Possible SN 1054 Petrograph by the Anasazi people (Chaco Canyon, New Mexico)

SN 1604

A New Star (Stella Nova), as bright and as red as Mars, was discovered on October 9 1604. At Padua the new star was observed by Galileo, while in Prague Kepler made careful observations and the Supernova now carries his name.



Kepler, *De Stella Nova in Pede Serpentarii,* (1606)

It is 400 years since a Supernova was last definitely observed within our own galaxy.

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Galileo gave a course of three lectures in Padua upon SN 1604 to a great audience. The main point he brought out concerning the new star was that it upset the received Aristotelian doctrine of the immutability of the heavens.

"Che ha a che fare la filosofia col misurare? Che importa al matematico se il cielo sia corruttibile e generabile? Se anche la nuova stella fosse di polenta, chi vieta ai matematici di osservarla e misurarla? Canchero, l'ha avuto torto questa stella a rovinare così la filosofia di costoro"

G. Galilei, Dialogo de Cecco Ronchitti da Bruzene in perpuosito de la Stella Nuova





First identification of SN



First modern identification by Baade & Zwicky (1934):

... at their maximum brightness they emit nearly as much light as the whole nebula in which they originate.

... their frequency is of the order of one super-nova per stellar system (nebula) per several centuries.

... Super-novae, initially, are quite ordinary stars whose masses are not greater than 10³³ gr. to 10³⁵ gr.

... The total energy emitted during the existence of the super-nova therefore is of the order of $\ldots 2.99 \times 10^{51}$ ergs.

... the phenomenon of a super-nova represents the transition of an ordinary star into a body of considerably smaller mass.

Thermonuclear vs. Core-Collapse Supernovae

Thermonuclear (Type Ia)

 Carbon-oxygen white dwarf (remnant of low-mass star)



 Accretes matter from companion

Core collapse (Type II, Ib/c)

- Degenerate iron core of evolved massive star
- Accretes matter by nuclear burning at its surface



Chandrasekhar limit is reached $-M_{Ch} \approx 1.5 M_{sun} (2Y_e)^2$ COLLAPSE SETS IN

Nuclear burning of C and O ignites → Nuclear deflagration ("Fusion bomb" triggered by collapse) Collapse to nuclear density Bounce & shock Implosion \rightarrow Explosion

Powered by nuclear binding energy

Powered by gravity

Comparable "visible" energy release of $\sim 3 \times 10^{51}$ erg

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Basic Spectral Classes

Type Ia seen in old stellar populations, while Ib, Ic & II seen in young populations \rightarrow core-collapse from massive stars

- Type Ia: No H, Si II, λ6355 seen
- Type Ib: No H, He I, λ5876 seen
- Type Ic: No H, no He or Si
- Type II: H seen (e.g. Balmer Hα λ6563, Hβ λ4861 etc)



Cox (2000)



Figure 1. Basic SN types spectra. A SN which near maximum (left panel) shows clear signature of H α is defined as type II, if it shows a strong SiII absorption at about 6150 Å is a type Ia, otherwise it is of type Ib/c (in the figure we show a SNIb characterized by strong He lines). Ten months later (right panel) SN Ia show strong emissions of [FeII] and [FeIII], SN Ib/c are dominated by [CaII] and [OI]. These same lines and strong H α emission are typical of SNII.

UNIVERSAL SUPERNOVA Ia LIGHT CURVE



LUMINOSITY DISTANCE

Distance estimates from SNe Ia light curves are derived from the luminosity distance

$$d_L = \left(\frac{L}{4\pi F}\right)^{\frac{1}{2}}$$

• F observed flux

It depends on the cosmological parameters

The luminosity distance is often expressed in terms of the magnitude

$$m = M + 5 \log_{10} \left(\frac{d_L}{Mpc}\right) + 25$$

M is the absolute magnitude (value at 10 pc)

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HUBBLE DIAGRAM





EXAMPLES OF SN II LIGHT CURVES



Core-Collapse SN Rate in the Milky Way



References: van den Bergh & McClure, ApJ 425 (1994) 205. Cappellaro & Turatto, astroph/0012455. Diehl et al., Nature 439 (2006) 45. Strom, Astron. Astrophys. 288 (1994) L1. Tammann et al., ApJ 92 (1994) 487. Alekseev et al., JETP 77 (1993) 339 and my update.

SN SUMMARY

Spectral Type	la	lb	lc	II
		No Hydrogen		
Spectrum	Silicon	No Silicon		
		Helium	No Helium	
Physical Mechanism	Nuclear explosion of low-mass star	Core collapse of evolved massive star (may have lost its hydrogen or even helium envelope during red-giant evolution)		
Light Curve	Reproducible	Large variations		
Neutrinos	Insignificant	\sim 100 $ imes$ Visible energy		
Compact Remnant	None	Neutron star (typically appears as pulsar) Sometimes black hole		
Rate / h ² SNu	$\textbf{0.36} \pm \textbf{0.11}$	0.14	± 0.07	$\textbf{0.71} \pm \textbf{0.34}$
Observed	Total \sim 5600 as of 2011 (Asiago SN Catalogue)			

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Stellar Collapse and Supernova Explosion



Stellar Collapse and Supernova Explosion





Gravitational binding energy:

 $E_b \approx 3 \times 10^{53} \text{ erg} \approx 17\% \text{ M}_{SUN} \text{ c}^2$

Energy sharing:

99%	Neutrinos
1%	Kinetic energy of explosion
	(1% of this into cosmic rays)
0.01%	Photons, outshine host galaxy

Neutrino luminosity:

 $\begin{array}{l} \mathsf{L}_{\mathsf{v}} \ \approx \ 3 \times 10^{53} \ \text{erg} \ / \ 3 \ \text{sec} \\ \ \approx \ 3 \times 10^{19} \ \mathsf{L}_{\mathsf{SUN}} \end{array}$

While it lasts, outshines the entire visible universe!

What determines the time scale?

Main neutrino reactions	Electron flavor $v_e + n \rightarrow p + e^ \overline{v}_e + p \rightarrow n + e^+$ Other flavors $v_e + N \rightarrow N + v$
Neutral-current scattering cross section	$\sigma(\nu N \rightarrow N\nu) = \frac{C_V^2 + 3C_A^2}{\pi} G_F^2 E_\nu^2 \approx 2 \times 10^{-40} \text{ cm}^2 \left(\frac{E_\nu}{100 \text{ MeV}}\right)^2$
Nucleon density	$n_{\rm B} = \frac{\rho_{\rm nuc}}{m_{\rm N}} \approx 1.8 \times 10^{38} {\rm cm}^{-3}$
Scattering rate	$\Gamma = \sigma n_{\rm B} \approx 1.1 \times 10^9 {\rm s}^{-1} \left(\frac{E_{\nu}}{100 {\rm MeV}}\right)^2$
Mean free path	$\lambda = (\sigma n_B)^{-1} \approx 28 cm \left(\frac{100 \text{ MeV}}{E_v}\right)^2$
Diffusion time	$t_{diff} \approx \frac{R^2}{\lambda} \approx 1.2 \sec\left(\frac{R}{10 \text{ km}}\right)^2 \left(\frac{E_v}{100 \text{ MeV}}\right)^2$

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Why No Prompt Explosion?

0.1 M_{sun} of iron has a nuclear binding energy ≈ 1.7 × 10⁵¹ erg
 Comparable to explosion energy

Dissociated Material (n, p, e, v)

zhock

Poissociat

- Shock wave forms within the iron core
- Dissipates its energy by dissociating the remaining layer of iron

ROLE OF NEUTRINOS

 Neutrinos produced in the hot, forming neutron star carry away the gravitational binding energy of the collapsing stellar core

$$E \approx 3 \times 10^{53} \left(\frac{M}{M_{sun}}\right)^2 \left(\frac{R}{10 \text{ km}}\right)^{-1} ergs$$

Neutrino energy $E_{\rm v}{\approx}$ 100 ${\times}$ E_{kin} of a SN explosion

 Neutrinos transfer energy to the collapsing stellar matter around the newly formed neutron star and could power the SN explosion

Characteristic SN energy unit: 10^{51} erg = 10^{44} J = 1 Bethe

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NEUTRINOS AND EXPLOSION MECHANISM

Paradigm: Explosions by the convectively supported neutrino-heating mechanism



- "Neutrino-heating mechanism": Neutrinos "revive" stalled shock by energy deposition [Colgate & White, 1966, Wilson, 1982, Bethe & Wilson, 1985]
- Convective processes & hydrodynamic instabilities enhance the heating mechanism [Herant et al. 1992, 1994; Burrows et al. 1995, Janka & Müller 1994, 1996; Fryer & Warren 2002, 2004; Blondin et al. 2003; Scheck et al. 2004,06,08]

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A CONVECTIVE ENGINE



Figure 5. A Convective Engine

(a) For simplicity, supernovae were often modeled in one dimension. A star was assumed to be spherically symmetric, its radius being the only spatial parameter that mattered. Doing simulations was therefore equivalent to doing physics in a long tube, even though the transfer of heat from one end of a pipe to the other is not very effective. (b) With the advent of multidimensional models, convection could occur. Hot, buoyant material could rise in one part of the star, to be replaced by cooler material failing from some other region. An in-out circuit is established that allows for the efficient and continuous transfer of heat out of the core and into the quasi-static layer. Energy from the gravitational collapse is thus converted into mechanical work as heat is being transferred between hot and cold reservoirs. In this sense, supernovae can be thought of as being powered by a simple convective engine.

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Neutrino-Driven Supernovae

- Stalled accretion shock still pushed outward to ~150km as matter piles up on the PNS, then recedes again
- Heating or gain region develops some tens of ms after bounce
- Convective overturn & shock oscillations "SASI" enhance the efficiency of v-heating, which finally revives the shock
- Big challenge: Show that this works!



Slide by B. Müller



Or with simpler schemes: e.g. IDSA+leakage Takiwaki et al. (2014)



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THREE PHASES OF NEUTRINO EMISSION

[Figure adapted from *Fischer et al. (Basel group), arXiv: 0908.1871*] 10. 8 M_{sun} progenitor mass

(spherically symmetric with Boltzmnann v transport)

Neutronization burst

• De-leptonization of outer

Shock breakout

core layers

Accretion

- Shock stalls ~ 150 km
- v powered by infalling matter

Cooling

 \bullet Cooling on ν diffusion time scale



MULTI-MESSENGER SIGNALS FROM SNE

[Nakamura, Horiuchi, Tamaka, Hayama, arXiV:1602.03028]



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GW FROM SNe

Millisecond bounce time reconstruction with nu trigger ICE-CUBE



FIG. 1: Typical Monte Carlo realization (red histogram) and reconstructed fit (blue line) for the benchmark case discussed in the text for a SN at 10 kpc.

[Halzen & Raffelt, arXiv:0908.2317]

See also [Pagliaroli, Vissani, Coccia & Fulgione arXiv:0903.1191]



SN 1987A NEUTRINOS

Sanduleak -69 202

Supernova 1987A 23 February 1987

<u>Neutrino Burst Observation :</u> First verification of stellar evolution mechanism



2002 Physics Nobel Prize

"for pioneering contributions to astrophysics, in particular for the detection of cosmic neutrinos"



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<u>Kamiokande-I to Kamiokande-II</u>

Kamiokande : Kamioka Nucleon Decay Experiment

3000-ton Imaging Water Cerenkov Detector





Today's backgrounds are tomorrow's signals.
Cherenkov Effect



IMB: 8000 ton Water Cherenkov Detector



-

E. N. Alexeyev (Hawaii, 2007)

The Baksan underground scintillation telescope



Liquid Scintillator Detector (LSD)

O. Saavedra (Hawaii, 2007)



NEUTRINO SIGNAL OF SN 1987A IN KAMIOKANDE



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NEUTRINO SIGNAL OF SUPERNOVA 1987A



Kamiokande-II (Japan) Water Cherenkov detector 2140 tons Clock uncertainty ±1 min

Irvine-Michigan-Brookhaven (US) Water Cherenkov detector 6800 tons Clock uncertainty ±50 ms

Baksan Scintillator Telescope (Soviet Union), 200 tons Random event cluster ~ 0.7/day Clock uncertainty +2/-54 s

Within clock uncertainties, signals are contemporaneous

SN 1987A Signal in LSD (Mont Blanc) ?



LSD (Liquid Scintillator Detector) in the Mont Blanc Tunnel (Oct. 1984 - March 1999) Supernova monitor for our galaxy 90 tons scintillator 200 tons iron (support structure)

- Observed a 5-event cluster
 4.72 hours before IMB/Kam-II
- Triggered autmatic SN alert
- Statistical fluctuation very unlikely
- No significant signal in IMB/Kam-II at LSD time
- No significant LSD signal at IMB time
- Interpretation as "double bang": Huge ν_e flux (~ 40 MeV) at LSD time
- LSD signal caused by interactions in iron of support structure
- Second bang ordinary multi-flavor signal

(Imshennik & Ryazhskaya, "A rotating collapsar and possible interpretation of the LSD neutrino signal from SN 1987A", astro-ph/0401613)

CROSS SECTION IN WC DETECTOR



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Angular Distribution of SN 1987A Neutrinos



Energy Distribution of SN 1987A Neutrinos



Kamiokande-II (Japan) Water Cherenkov detector 2140 tons

Irvine-Michigan-Brookhaven (US) Water Cherenkov detector 6800 tons

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INTERPRETING SN 1987A NEUTRINOS

[e.g.,B. Jegerlehner, F. Neubig and G. Raffelt, PRD **54**, 1194 (1996); <u>A.M.</u>, and G. Raffelt, PRD **72**, 063001 (2005)]



In agreement with the most recent theoretical predictions (i.e. Basel & Garching models)

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SN NEUTRINO LIGHT CURVE FROM SN 1987A

[Loredo & Lamb, astro-ph/0107260 ; Pagliaroli, Vissani, Costantini & Ianni,arXiV:0810.0466]

Figure 6: Pagliaroli et al. model: antineutrino luminosity and average energy in the best fit point.



 $R_c = 16 \text{ km}, \quad T_c = 4.6 \text{ MeV}, \quad \tau_c = 4.7 \text{ s}, \text{ cooling}$ $M_a = 0.2 M_{\odot}, \quad T_a = 2.4 \text{ MeV}, \quad \tau_a = 0.6 \text{ s}.$ accretion

Light curve in reasonable agreement with generic expectations of delayed explosion scenario

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SN NEUTRINO SPECTRUM FROM SN1987A

[Yuksel & Beacom, astro-ph/0702613]



Original SN ν energy spectra expected to be quasi-thermal

SN1987A inferred v energy spectrum shows strong deviations from quasi-thermal distribution:

Possible effects of:

- neutrino mixing
- v-v interactions
- v decay
- nonstandard v interactions
- additional channels of energy exchange among flavors

Possible to reconcile detection and theory... Still many open questions!

Based on the handful of SN neutrinos which were detected that day, approximatively one theory paper has published every ten days....



... for the last thirty years!

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PARTICLE PHYSICS BOUNDS FROM SN 1987A



- Exotic neutrino properties
 - Axion-like particles
- Energy-loss and novel particles

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BOUND ON SECRET NEUTRINO INTERACTIONS

$$L = g\phi v\overline{v}$$

 ϕ new scalar mediator with mass M

Four fermion approximation $G = \frac{1}{\sqrt{4\pi}} \frac{g^2}{M^2}$

Requiring that v from cosmic sources travel through the CvB without scattering induced by the secret interactions leads to upper limits on the new coupling.



SN 1987A bound $G \leq 10^{-8} GeV^{-2}$ [Kolb & Turner, PRD 36, 2895 (1987)]

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SN 1987A BOUNDS ON NEUTRINO VELOCITY



SN 1987A few events provide the most stringent constraints on v velocity. Crucial for comparison with recent OPERA claim



Table 1. Superluminal Neutrino Velocity Observations and Bounds [Evslin, 1111.0733]

OPERA	2009-2011							
Energy	Neutrinos	(v-c)/c						
10-50 GeV	16,111 ν's (97% ν _μ '2)	2.48 ± 0.28 (stat.) ± 0.30 (syst.) $\times 10^{-5}$						
Distance:	ance: 730 km from CNGS (CERN) to OPERA (Gran Sasso)							
MINOS	May 2005-February 2006							
Energy: 3 GeV	Neutrinos	(v-c)/c						
(tail to 120 GeV)	473 ν's (93% ν _μ 's)	$5.1 \pm 1.3 \text{ (stat.)} \pm 2.6 \text{ (sys.)} \times 10^{-5}$						
Distance:	Distance: 734 km: Near Detector (FermiLab) to Soudan iron mine							
Kamiokande II	7:35 UT, February 23rd, 1987							
Energy	Neutrinos	ν 's \subset 13 sec., \lesssim 3 hrs before γ 's,						
7.5-36 MeV	$12 p_e$'s	$(v - c)/c < 3 \times 10^{-9}$ or 2×10^{-12}						
Distance:	160,000 lys: Tarantula Nebula to Kamioka Observatory							
Irvine-Michigan-Brookhaven	7:35 UT, February 23rd, 1987							
Energy	Neutrinos	ν 's \subset 6 sec., \lesssim 3 hrs before γ 's,						
20-40 MeV	8 Ve's	$(v-c)/c < 3 \times 10^{-9} \text{ or } 2 \times 10^{-12}$						
Distance	160,000 lys: Tarantula Nebula to Morton-Thiokol salt mine							

AXION-LIKE PARTICLES (ALPs)

$$L_{a\gamma} = -\frac{1}{4} g_{a\gamma} F_{\mu\nu} \widetilde{F}_{\mu\nu} a = g_{a\gamma} \vec{E} \cdot \vec{B} a$$

Primakoff process: Photon-ALP transitions in external static E or B field

Photon-ALP conversions in macroscopic B-fields

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ALPS CONVERSIONS FOR SN 1987A

Milky-Way

[Brockway, Carlson, Raffelt, astro-ph/9605197, Masso and Toldra, astro-ph/9606028]

SN 1987A



ALPs produced in SN core by Primakoff process

ALP-photon conversions in the Galactic B-fields

No excess gammarays in coincidence with SN 1987A

SMM Satellite

In [Payez, Evoli, Fischer, Giannotti, <u>A.M.</u> & Ringwald, 1410.3747] we revaluate the bound with

- state-of-art models for SNe and Galactic B-fields
- accurate microscopic description of the SN plasma

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ALP-PHOTON FLUXES FOR SN 1987A

[Payez, Evoli, Fischer, Giannotti, <u>A.M.</u> & Ringwald, 1410.3747]



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GAMMA-RAY OBSERVATION FROM SMM SATELLITE





NEW BOUND ON ALPs FROM SN 1987A

[Payez, Evoli, Fischer, Giannotti, <u>A.M.</u> & Ringwald, 1410.3747]



 $g_{a\gamma} \le 5.3 \times 10^{-12} \ GeV^{-1}$ for $m_a < 4.4 \times 10^{-10} \text{eV}$

SN 1987A provides the strongest bound on ALP-photon coversions for ultralight ALPs

ENERGY-LOSS ARGUMENT



Assuming that the SN 1987A neutrino burst was not shortened by more than $\sim \frac{1}{2}$ leads to an approximate requirement on a novel energy-loss rate of

$$\epsilon_{\chi} < 10^{19} \text{ erg g}^{-1} \text{ s}^{-1}$$

for
$$\rho \approx 3 \times 10^{14}$$
 g cm⁻³ and T ≈ 30 MeV
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AXION EMISSION FROM A NUCLEAR MEDIUM



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SN 1987A AXION LIMITS



Hadronic axion ($m_a \sim 1 \text{ eV}$, $f_a \sim 10^6 \text{ GeV}$) not excluded by SN 1987A. Possible hot-dark matter candidate. The "hadronic axion window" is closed by cosmological mass bounds.

SN 1987A BOUND ON HIDDEN PHOTONS



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SN 1987A BOUND ON KeV STERILE NEUTRINOS



KeV sterile v are produced in a SN core by the mixing with active v.

• For sufficiently small mixing θ , v_s escape the core immediately after the production contributing to the energy-loss.

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EUTURE SNUR UTRANS OBS SNULLONE

NEUTRINO DETECTION METHODS



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Milky-Way SN

Excellent statistics (10⁴ events for 10 kpc) High-sensitivity to explosion scenario 1 SN ~ 40 years

SNe in nearby galaxies

Few to 10 neutrinos per SN, but requires a Mtonclass detector 1 SN ~ year

Diffuse Supernova Neutrino Background (DSNB)

Neutrinos from all past core-collapse SNe; emission is averaged, no timing or direction (faint) signal is always there

PROBABILITY OF MILKY-WAY SN

[Kistler, Yuksel, Ando, Beacom & Suzuki, 0810.1959]



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GALACTIC SUPERNOVA DISTANCE DISTRIBUTION

[<u>A.M.</u>, Raffelt, Serpico, astro-ph/ 0604300]



Average distance 10.7 kpc, rms dispersion 4.9 kpc (11.9 kpc and 6.0 kpc for SN Ia distribution)

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Large Detectors for Supernova Neutrinos



In brackets events for a "fiducial SN" at distance 10 kpc

LARGE DETECTORS FOR SN NEUTRINOS

[A.M., Tamborra, Janka, Saviano, Scholberg et al., arXiv:1508.00785 [astro-ph.HE]]

Detector	Type	Mass (kt)	Location	Events	Flavors	Status
Super-Kamiokande	H_2O	32	Japan	7,000	$\bar{ u}_e$	Running
LVD	$C_n H_{2n}$	1	Italy	300	$\bar{\nu}_e$	Running
KamLAND	$C_n H_{2n}$	1	Japan	300	$\bar{\nu}_e$	Running
Borexino	$C_n H_{2n}$	0.3	Italy	100	$\bar{ u}_e$	Running
IceCube	Long string	(600)	South Pole	(10^{6})	$\bar{ u}_e$	Running
Baksan	$C_n H_{2n}$	0.33	Russia	50	$\bar{ u}_e$	Running
$MiniBooNE^*$	$C_n H_{2n}$	0.7	USA	200	$\bar{\nu}_e$	(Running)
HALO	\mathbf{Pb}	0.08	Canada	30	$ u_e, u_x$	Running
Daya Bay	$C_n H_{2n}$	0.33	China	100	$\bar{ u}_e$	Running
$NO\nu A^*$	$C_n H_{2n}$	15	USA	4,000	$\bar{\nu}_e$	Turning on
SNO+	$C_n H_{2n}$	0.8	Canada	300	$\bar{\nu}_e$	Near future
MicroBooNE*	Ar	0.17	USA	17	ν_e	Near future
DUNE	Ar	34	USA	3,000	ν_e	Proposed
Hyper-Kamiokande	H_2O	560	Japan	110,000	$\bar{\nu}_e$	Proposed
JUNO	$C_n H_{2n}$	20	China	6000	$\bar{\nu}_e$	Proposed
RENO-50	$C_n H_{2n}$	18	Korea	5400	$\bar{\nu}_e$	Proposed
PINGU	Long string	(600)	South Pole	(10^{6})	$ar{ u}_e$	Proposed

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SUPERNOVA EARLY WARNING SYSTEM (SNEWS)

[P.Antonioli et al., astro-ph/0406214]

Neutrinos several hours before light



Neutrino observations can alert astronomers several hours in advance to a SN. To avoid false alarms, require alarm from at least two experiments



SUPER-KAMIOKANDE DETECTOR



SK is a cylindrical tank containing 50000 ton of light water surrounded by photomultipliers, located underground in the Kamioka mine in Japan.

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Simulated Supernova Signal at Super-Kamiokande



[Totani, Sato, Dalhed & Wilson, ApJ 496 (1998) 216]
ICECUBE NEUTRINO TELESCOPE AT SOUTH POLE



SN NU SIGNAL IN ICECUBE

[A.M., Tamborra, Janka, Saviano, Scholberg et al., arXiv:1508.00785 [astro-ph.HE]]



High statistics reconstruction of the nu light curve. Possible to distinguish the different post-bounce phases.

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MILLISECOND BOUNCE TIME RECONSTRUCTION



ICECUBE

FIG. 1: Typical Monte Carlo realization (red histogram) and reconstructed fit (blue line) for the benchmark case discussed in the text for a SN at 10 kpc.

30

[Halzen & Raffelt, arXiv:0908.2317]

Possible also in Super-K [see Pagliaroli, Vissani, Coccia & Fulgione arXiv:0903.1191]



Tamborra, Hanke, Müller, Janka & Raffelt, arXiv:1307.7936

Necessary high statistics and high time resolution

Convective motions lead to large-amplitude oscillations of the stalled shock with a period of ~ 10 ms



Icecube is ok!

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NEXT-GENERATION DETECTORS



SN NU SIGNAL IN FUTURE DETECTORS

[A.M., Tamborra, Janka, Saviano, Scholberg et al., arXiv:1508.00785 [astro-ph.HE]]



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SN NU SIGNAL IN DM DETECTORS

[Lang,McCabe, Reichard, Selvi & Tamborra, arXiv:1606.09243 [astro-ph.HE]]



DARWIN will be able to clearly reconstruct the neutrino light-curve and to differentiate among phases of neutrino signal. Partial sensitivity with XENONnT/LZ.

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A REAPPRISAL OF AXION EMISSION WITH STATE OF ART SIMULATIONS

[Fischer, Chakraborty, Giannotti A.M., Payez & Ringwald, 1605.08780]

18 M_{sun} progenitor mass

(spherically symmetric with Boltzmnann v transport)





@ Super-Kamiokande



NEUTRINOS FROM ALL COSMEC SUPERNAVAE

Local Group of Galaxies



OBSERVED SUPERNOVAE IN THE LOCAL UNIVERSE

[Kistler, Yuksel, Ando, Beacom & Suzuki, 0810.1959]



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Detection of Neutrinos from Supernovae in Nearby Galaxies

[S. Ando, J. Beacom, and Y. Yuksel, astro-ph/0503321]

Mton Cherenkov

Cumulative SN rate



Reconstruction of SN neutrino spectrum by the patient accumulation of ~1 neutrino per supernova from galaxies within 10 Mpc, in which one expects at least 1(2) SN per year.

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NEUTRINO EVENTS RATE FROM EXTRAGALACTIC SNe

[Kistler et al., 0810:1959]

TABLE I: Approximate neutrino event yields for core-collapse supernovae from representative distances and galaxies, as seen in various detectors with assumed fiducial volumes. Super-Kamiokande is operating, and Hyper-Kamiokande and Deep-TITAND are proposed.

		32 kton	0.5 Mton	5 Mton	
		(SK)	(HK)	(Deep-TITAND)	
10 kpc	(Milky Way)	10^{4}	10^{5}	10^{6}	
1 Mpc	(M31, M33)	1	10	10^{2}	
3 Mpc	(M81, M82)	10^{-1}	1	10	

TABLE II: Core-collapse supernova candidates from 1999-2008 within 6 Mpc, with their expected neutrino event yields (E_{e^+} > 18 MeV) in a 5 Mton detector.

SN	Туре	Host	D [Mpc]	ν events
2002hh	II-P	NGC 6946	5.6	2.4
2002kg	IIn/LBV	NGC 2403	3.3	6.8
2004am	II-P	NGC 3034 (M82)	3.53	5.9
2004dj	II-P	NGC 2403	3.3	6.8
2004et	II-P	NGC 6946	5.6	2.4
2005af	II-P	NGC 4945	3.6	5.7
2008S	IIn	NGC 6946	5.6	2.4
2008bk	II-P	NGC 7793	3.91	4.8
2008iz	II?	NGC 3034 (M82)	3.53	5.9
NGC 300-T	II?	NGC 300	2.15	16.0

Supernova Explosion in M82: Exciting, but No Neutrinos



The M82 galaxy before (top) and after (bottom) its new supernova on Jan. 22 (Photo: UCL/University of London Observatory/Steve Fossey/Ben Cooke/Guy Pollack/Matthew Wilde/Thomas Wright)

22 Jan. 2014

By Erin Weeks

In the early morning hours of January 22, the Earth turned spectator to a celestial event the likes of which hadn't been seen in nearly three decades. The explosive death of a white dwarf star in Messier 82 (M82), a nearby galaxy, quickly ignited the astronomy world.

The supernova is exciting for a number of reasons that other outlets have well outlined — but unfortunately for Kate Scholberg, neutrinos are not one of them. Scholberg, a Duke University physics professor, studies the mysterious, nearly-massless particles at Super-K, a detector located deep in the mountains of Japan. Super-K was designed to spot neutrinos as they speed through Earth, revealing information about their sources, which can include the sun, cosmic rays, and supernovae.

"M82 is too far away for us to see any neutrinos from it," Scholberg wrote in an email. "It's about 11.4 million light years from us, meaning that the chance of seeing even a single neutrino from a core-collapse supernova in current detectors is probably a few percent or less (of course, we'll look)." A galactic SN explosion is a spectacular event which will produce an enormous number of detectable v, but it is a <u>rare</u> event (~ 3/century) ...



 \dots Conversly, there is a guaranteed v background produced by all the past Supernovae in the Universe, but leading to much less detectable events.

THE DIFFUSE SUPERNOVA NEUTRINO BACKGROUND (DSNB)





WHAT CAN WE LEARN FROM DSNB?

In principle, we can extract information on:

- Star formation rate
- Neutrino masses and mixing parameters
- SN neutrino energies

... not all at the same time, however!

(degeneracy of effects)

BACKGROUND IN SK FOR DSNB SIGNAL



Below ~15–20 MeV, bkgd dominated by spallation products (made by atmospheric μ) and by reactor $\overline{\nu}_e$.

For $E_{\rm v} \in$ [20-30] MeV, the bkg of low-energy atmospheric $\overline{\rm v}_e$ is relatively small.

But, in this window, there is a large background due to "invisible" μ (i.e. below Cherenkov emission threshold) decay products, induced by low energy atmospheric v_{μ} and \overline{v}_{μ} .

DSNB signal should manifest as distortion of the bkg spectra.

No distortion \longrightarrow flux limit

Alessandro Mírízzí



DSNB signal should manifest as distortion of the bkg spectra.

No distortion → flux limit

FIG. 11. Spectra of the four remaining backgrounds in the signal Cherenkov angle region with all reduction cuts applied. The ν_{μ} CC channel is from decay electron data; the other three are from MC. All are scaled to the SK-I LMA best fit result.

Super-Kamiokande collaboration recently investigated the DSNB flux using 2853 days of data [Bays et al., arXiV:1111:5031]. It fixed an upper bound on DSNB signal:

$$J_{\overline{v_e}} \le 3 \text{ cm}^{-2} \text{ s}^{-1}$$

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SN NEUTRINO EMISSION LIMIT FROM DSNB

[Bays et al., arXiV:1111.5031, see also Vissani & Pagliaroli, 1102.0447]



The SK limit is close to the most recent theoretical predictions ... but Super-K is background limited.

SK-Gd Project



 $v_{\rm e}$ can be identified by delayed coincidence.

[reaction schematic by M. Nakahata]

[G.L.Fogli, E.Lisi, A.M., and D.Montanino, hep-ph/0412046]

Mton Cherenkov



Below ~15–20 MeV, bkgd dominated by spallation products (made by atmospheric μ) and by reactor \overline{v}_e .

For $E_{\rm v} \in$ [20-30] MeV, the bkg of low-energy atmospheric $\overline{\rm v}_e$ is relatively small.

But, in this window, there is a large background due to "invisible" μ (i.e. below Cherenkov emission threshold) decay products, induced by low energy atmospheric ν_{μ} and $\overline{\nu}_{\mu}$.

Adding Gd [J.F.Beacom, and M.R.Vagins, hep-ph/0309300], spallation ~eliminated, invisible μ reduced by ~5. The analysis threshold lowered.

A few clean events/y in Super-K with Gd ! In HK after 1 year, the DSNB signal detectable at 6σ level

Alessandro Mírízzí

EGADS

Evaluating Gadolinium's Action on Detector Systems

- To study the Gd water quality with actual detector materials.
- The detector fully mimic Super-K detector;

SUS frame, PMT and PMT case, black sheets, etc.

Tests for Hyper-K; 13 HPDs





Hiroyuki Sekiya

NEUTRINO2016 London

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SK & T2K Joint Statement on "SK-Gd" Jan.30, 2016

On June 27, 2015, the Super-Kamiokande collaboration approved the SK-Gd project which will enhance neutrino detectability by dissolving gadolinium in the Super-K water.

T2K and SK will jointly develop a protocol to make the decision about when to trigger the SK-Gd project, taking into account the needs of both experiments, including preparation for the refurbishment of the SK tank and readiness of the SK-Gd project, and the T2K schedule including the J-PARC MR power upgrade. Given the currently anticipated schedules, the expected time of the refurbishment is 2018.

After years of testing and study – culminating in these powerful EGADS results – no technical showstoppers have been encountered. And so...

June 27, 2015: The Super-Kamiokande Collaboration approved the addition of gadolinium to the detector, pending discussions with T2K.



January 30, 2016: The T2K Collaboration approved addition of gadolinium to Super-Kamiokande, with the precise timing to be jointly determined based on the needs of both projects.



July 26, 2017: The official start time of draining the SK tank to prepare for Gd loading is decided to be <u>June 1, 2018</u>.



Slide from Mark Vagins

Hall G being filled with equipment for the gadolinium loading of Super-Kamiokande; January 30th, 2017

Expected timeline for SK-Gd



<u>We should have collected some</u> <u>new supernova neutrinos for the</u> <u>theorists to study with within</u> <u>three years from today!</u>

DSNB IN LARGE FUTURE DETECTORS

[<u>A.M.</u>, Tamborra, Janka, Saviano, Scholberg et al., arXiv:1508.00785 [astro-ph.HE]]



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CONSTRAINT OF NU INVISIBLE DECAY FROM DSNB

[Fogli, Lisi, <u>A.M.</u>, Montanino, hep-ph/0401227]



Nu decay in Majoron

$$\nu \rightarrow \nu' + \phi$$

DSNB can probe lifetimes of cosmological interest

 $\frac{\tau_i E}{m_i} \le 1 / H_0$



DSNB spectrum larger, comparable or smaller than the standard one



VACUUM OSCILLATIONS

• Two flavor mixing
$$\begin{pmatrix} v_e \\ v_\mu \end{pmatrix} = \begin{pmatrix} \cos\theta & \sin\theta \\ -\sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix}$$

Each mass eigenstates propagates as e^{ipz} with $p_i = \sqrt{E^2 - m^2} \approx E - \frac{m_i^2}{2E}$

$$\left| \boldsymbol{v}_{\mu}(z) \right\rangle = -\sin\theta \ e^{-ip_{1}z} \left| \boldsymbol{v}_{1} \right\rangle + \cos\theta \ e^{-ip_{2}z} \left| \boldsymbol{v}_{2} \right\rangle$$

2 v oscillation probability $P(\boldsymbol{v}_{e} \rightarrow \boldsymbol{v}_{\mu}) = \left| \left\langle \boldsymbol{v}_{\mu}(z) \left| \boldsymbol{v}_{e}(0) \right\rangle \right|^{2} = \sin^{2} 2\theta \sin^{2} \left(\frac{\Delta m^{2}L}{4E} \right)$



Bruno Pontecorvo (1967)



Dubna 1988. Neutrino oscillations.

SN v FLAVOR TRANSITIONS

The flavor evolution in matter is described by the non-linear MSW equations:

$$i\frac{d}{dx}\psi_{v} = \left(H_{vac} + H_{e} + H_{vv}\right)\psi_{v}$$

In the standard 3v framework

•
$$H_{vac} = \frac{U M^2 U^{\dagger}}{2E}$$

• $H_e = \sqrt{2}G_F \operatorname{diag}(N_e, 0, 0)$
• $H_{vv} = \sqrt{2}G_F \int (1 - \cos \theta_{pq}) \left(\rho_q - \overline{\rho}_q\right) dq$

Kinematical mass-mixing term

Dynamical MSW term (in matter)

Neutrino-neutrino interactions term (non-linear)

3v FRAMEWORK

• Mixing parameters: $U = U(\theta_{12}, \theta_{13}, \theta_{23}, \delta)$ as for CKM matrix



 $\textbf{c}_{12}\text{=}\cos\,\theta_{12},\,\text{etc.},\,\delta$ CP phase

Mass-gap parameters:

arameters:
$$M^2 = \left(\begin{array}{c} -\frac{\delta m^2}{2}, +\frac{\delta m^2}{2}, \pm \Delta m^2\right)$$

"solar" "atmospheric"
 $V_3 + \Delta m^2$ inverted hierarchy
 $V_1 + \delta m^{2/2}$ $V_1 + \delta m^{2/2}$
 $V_2 - \delta m^{2/2}$ $V_2 - \delta m^{2/2}$

normal hierarchy $v_3 - \Delta m^2$

SN neutrinos are sensitive to the unknown mass hierarchy

STATUS OF NEUTRINO OSCILLATIONS



GLOBAL OSCILLATION ANALYSIS (2017)

[Capozzi et al., arXiv:1703.0447]





Collective flavor transitions at low-radii [O (10² - 10³ km)]

Alessandro Mírízzí
Collective Supernova Nu Oscillations since 2006

Two seminal papers in 2006 triggered a torrent of activities Duan, Fuller, Qian, astro-ph/0511275, Duan et al. astro-ph/0606616

Balantekin, Gava & Volpe [0710.3112]. Balantekin & Pehlivan [astro-ph/0607527]. Blennow, Mirizzi & Serpico [0810.2297]. Cherry, Fuller, Carlson, Duan & Qian [1006.2175, 1108.4064]. Cherry, Wu, Fuller, Carlson, Duan & Qian [1109.5195]. Cherry, Carlson, Friedland, Fuller & Vlasenko [1203.1607]. Chakraborty, Choubey, Dasgupta & Kar [0805.3131]. Chakraborty, Fischer, Mirizzi, Saviano, Tomàs [1104.4031, 1105.1130]. Choubey, Dasgupta, Dighe & Mirizzi [1008.0308]. Dasgupta & Dighe [0712.3798]. Dasgupta, Dighe & Mirizzi [0802.1481]. Dasgupta, Dighe, Raffelt & Smirnov [0904.3542]. Dasgupta, Dighe, Mirizzi & Raffelt [0801.1660, 0805.3300]. Dasgupta, Mirizzi, Tamborra & Tomàs [1002.2943]. Dasgupta, Raffelt & Tamborra [1001.5396]. Dasgupta, O'Connor & Ott [1106.1167]. Duan, Fuller, Carlson & Qian [astroph/0608050, 0703776, 0707.0290, 0710.1271]. Duan, Fuller & Qian [0706.4293, 0801.1363, 0808.2046, 1001.2799]. Duan, Fuller & Carlson [0803.3650]. Duan & Kneller [0904.0974]. Duan & Friedland [1006.2359]. Duan, Friedland, McLaughlin & Surman [1012.0532]. Esteban-Pretel, Mirizzi, Pastor, Tomàs, Raffelt, Serpico & Sigl [0807.0659]. Esteban-Pretel, Pastor, Tomàs, Raffelt & Sigl [0706.2498, 0712.1137]. Fogli, Lisi, Marrone & Mirizzi [0707.1998]. Fogli, Lisi, Marrone & Tamborra [0812.3031]. Friedland [1001.0996]. Gava & Jean-Louis [0907.3947]. Gava & Volpe [0807.3418]. Galais, Kneller & Volpe [1102.1471]. Galais & Volpe [1103.5302]. Gava, Kneller, Volpe & McLaughlin [0902.0317]. Hannestad, Raffelt, Sigl & Wong [astro-ph/0608695]. Wei Liao [0904.0075, 0904.2855]. Lunardini, Müller & Janka [0712.3000]. Mirizzi, Pozzorini, Raffelt & Serpico [0907.3674]. Mirizzi & Serpico [1111.4483]. Mirizzi & Tomàs [1012.1339]. Pehlivan, Balantekin, Kajino & Yoshida [1105.1182]. Pejcha, Dasgupta & Thompson [1106.5718]. Raffelt [0810.1407, 1103.2891]. Raffelt & Sigl [hep-ph/0701182]. Raffelt & Smirnov [0705.1830, 0709.4641]. Raffelt & Tamborra [1006.0002]. Sawyer [hep-ph/0408265, 0503013, 0803.4319, 1011.4585]. Sarikas, Raffelt, Hüdepohl & Janka [1109.3601]. Sarikas, Tamborra, Raffelt, Hüdepohl & Janka [1204.0971]. Saviano, Chakraborty, Fischer, Mirizzi [1203.1484]. Wu & Qian [1105.2068].....

NEUTRINO-NEUTRINO HAMILTONIAN

$$i\frac{d}{dx}\psi_{v}^{(i)} = \left(H_{vac} + H_{e} + \sum_{j}H_{vv}^{(ij)}\right)\psi_{v}^{(i)}$$

In early studies the neutrino-neutrino Hamiltonian was assumed diagonal in flavor basis: No contribution to flavor evolution!

Critical examination of this assumption by J.Pantaleone [PLB 287, 128 (1992)]

Low-energy neutral current Hamiltonian for v-v interactions possesses an U(N) symmetry. A diagonal H_{vv} doesn't respect this symmetry.

Pantaleone proposed a modified form of $H_{_{\!\rm VV}}$ which contains non-zero off-diagonal terms

It respects U(N)

symmetry

$$H_{\nu\nu} = A \begin{pmatrix} \left| \nu_{e} \right|^{2} & \nu_{e} \nu_{\mu}^{*} \\ \nu_{e}^{*} \nu_{\mu} & \left| \nu_{\mu} \right|^{2} \end{pmatrix}$$
$$A = \sqrt{2}G_{F} \left(1 - \cos \theta_{ij} \right)$$

NEUTRINO FLAVOR CONVERSIONS IN A NEUTRINO BACKGROUND

Since H_{vv} cannot change the *total* flavor of the system, v-v interactions do contribute to the flavor evolution only when the "propagating" and "background" neutrinos do exchange momenta



If all the v in the bkg are in the same flavor state $v_x = \cos \alpha v_e + \sin \alpha v_\mu$

$$H_{vv} \propto \frac{\sqrt{2}G_F n_2}{2} \begin{pmatrix} \cos 2\alpha & \sin 2\alpha \\ \sin 2\alpha & -\cos 2\alpha \end{pmatrix} \implies P(v_e \to v_\mu) \approx \sin^2 2\alpha (G_F n_2 L)^2 / 2$$

[Friedland & Lunardini, hep-ph/0304055]

However, one cannot distinguish btw beam and bkg. Instrinsic non-linear problem !

SELF-INDUCED SPECTRAL SPLITS

[Fogli, Lisi, Marrone, <u>A.M.</u>, arXiV: 0707.1998 [hep-ph], Duan, Carlson, Fuller, Qian, astro-ph/0703776, Raffelt and Smirnov, 0705.1830 [hep-ph], Dasgupta, Dighe, Raffelt & Smirnov, arXiv:0904.3542 [hep-ph], Duan & Friedland, arXiv: 1006.2359, <u>A.M.</u> & Tomas, arXiv:1012.1339, Choubey, Dasgupta, Dighe, <u>A.M.</u>, 1008.0308....]



Strong dependence of collective oscillations on mass hierarchy and on the energy ("splits")

Splits possible in both normal and inverted hierarchy, for $v \& \overline{v}$!! Alessandro Mirizzi Karpacz, 26 February 2019

SPONTANEOUS SYMMETRY BREAKING IN SELF-INDUCED OSCILLATIONS

- Symmetries have been used to reduce the complexity of the SN v flavor evolution (e.g. the bulb model).
- However, v can lead to a spontaneous symmetry breaking (SSB) of the symmetry inherent to the initial conditions [Raffelt, Sarikas, Seixas, 1305.7140].
- Small deviations from the space/time symmetries of the bulb model have to be expected. Can these act as seed for new instabilities?

FIRST INVESTIGATIONS WITH TOY MODELS

- With a simple toy model in [Mangano, <u>A.M.</u> & Saviano, 1403.1892] it has been shown that self-interacting v can break translational symmetries in space and time.
- By a stability analysis in [Duan & Shalgar, 1412.7097] is has been found that self-interacting v can break the spatial symmetries of a 2D model.

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2D MODEL FOR SELF-INTERACTING ν

[Duan & Shalgar, 1412.7097]



Nu evolving in the plane (x,z)emitted from an infinite boundary at z=0, in only two directions (L and R). Excess of v_e over $\overline{v_e}$.



<u>A. M.</u>, Mangano and Saviano, [arXiv:1503.03485 [hep-ph]].



Large variations in the x direction at smaller and smaller scales.

Planes of common phase broken.

Coherent behavior of oscillation lost.

RECENT PAPERS ON SSB

- G. Mangano, <u>A. M.</u> and N. Saviano, "Damping the neutrino flavor pendulum by breaking homogeneity," Phys. Rev. D 89, no. 7, 073017 (2014) [arXiv:1403.1892 [hep-ph]].
- H. Duan and S. Shalgar, "Flavor instabilities in the neutrino line model," Phys. Lett. B 747, 139 (2015) [arXiv:1412.7097 [hep-ph]].
- <u>A. M.</u>, G. Mangano and N. Saviano, "Self-induced flavor instabilities of a dense neutrino stream in a two-dimensional model," Phys. Rev. D 92, no. 2, 021702 (2015) [arXiv:1503.03485 [hep-ph]].
- S. Chakraborty, R. S. Hansen, I. Izaguirre and G. Raffelt, "Self-induced flavor conversion of supernova neutrinos on small scales," JCAP 1601, no. 01, 028 (2016) [arXiv:1507.07569 [hep-ph]].
- S. Abbar and H. Duan, "Neutrino flavor instabilities in a time-dependent supernova model," Phys. Lett. B 751, 43 (2015) [arXiv:1509.01538 [astro-ph.HE]].
- B. Dasgupta and <u>A. M.</u>, "Temporal Instability Enables Neutrino Flavor Conversions Deep Inside Supernovae," Phys. Rev. D 92, no. 12, 125030 (2015) [arXiv:1509.03171 [hep-ph]].

F. Capozzi, B. Dasgupta and <u>A. M.</u>, "Self-induced temporal instability from a neutrino antenna," JCAP 1604, no. 04, 043 (2016) [arXiv:1603.03288 [hep-ph]].
Alessandro Mírízzi

FLAVOR CONVERSIONS NEAR SN CORE?



Most of the studies assume no flavor conversion at r < 50 km (only synchronized oscillations). After selfinduced conversions develop with a rate ~ $\int \omega \mu$ [see, e.g., Hannestad et al, astro-ph/0608695]





However, since more than a decade Ray Sawyer is pointing out that close to nu-sphere nu angular distributions of different species are rather different. This would lead to a new flavor instability (absent assuming equal angular distributions). The outcome would be a possible complete flavor mixing of the outgoing stream just above the nu-sphere. Fast rate $\sim \mu$

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WHY SHOULD WE WORRY ABOUT THESE EFFECTS?

• If flavor changes occur in the deepest SN regions, they would modify the neutrino heating behind the stalled shock wave, possibly helping a SN to explode.



Exploding. Spectral swap (by hands) behind shock front [Suwa et al., 1106.5487]

• If flavor equilibrium would occur close to the nu-sphere, all further oscillation effects (self-indued, matter...) would be washed-out. Crucial to predict observable SN nu signal.

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FAST FLAVOR CONVERSIONS NEAR SN CORE

PHYSICAL REVIEW D 72, 045003 (2005)

Speed-up of neutrino transformations in a supernova environment

R.F. Sawyer

Department of Physics, University of California at Santa Barbara, Santa Barbara, California 93106, USA (Received 8 April 2005; published 5 August 2005)

When the neutral current neutrino-neutrino interaction is treated completely, rather than as an interaction among angle-averaged distributions, or as a set of flavor-diagonal effective potentials, the result can be flavor mixing at a speed orders of magnitude faster than that one would anticipate from the measured neutrino oscillation parameters. It is possible that the energy spectra of the three active species of neutrinos emerging from a supernova are nearly identical.

PHYSICAL REVIEW D 79, 105003 (2009)

Multiangle instability in dense neutrino systems

R. F. Sawyer

Department of Physics, University of California at Santa Barbara, Santa Barbara, California 93106, USA (Received 18 April 2008; published 6 May 2009)

We calculate rates of flavor exchange within clouds of neutrinos interacting with each other through the standard model coupling, assuming a conventional mass matrix. For cases in which there is an angular dependence in the relation among intensity, flavor, and spectrum, we find instabilities in the evolution equations and greatly speeded-up flavor exchange. The instabilities are categorized by examining linear perturbations to simple solutions, and their effects are exhibited in complete numerical solutions to the system. The application is to the region just under the neutrino surfaces in the supernova core.

PRL 116, 081101 (2016) PHYSICAL REVIEW LETTERS week ending 26 FEBRUARY 2016

Neutrino Cloud Instabilities Just above the Neutrino Sphere of a Supernova

R. F. Sawyer

Department of Physics, University of California at Santa Barbara, Santa Barbara, California 93106, USA (Received 7 September 2015; revised manuscript received 2 January 2016; published 25 February 2016)

Most treatments of neutrino flavor evolution, above a surface of the last scattering, take identical angular distributions on this surface for the different initial (unmixed) flavors, and for particles and antiparticles. Differences in these distributions must be present, as a result of the species-dependent scattering cross sections lower in the star. These lead to a new set of nonlinear equations, unstable even at the initial surface with respect to perturbations that break all-over spherical symmetry. There could be important consequences for explosion dynamics as well as for the neutrino pulse in the outer regions.

Literature on Fast Flavor Conversion

- 1. Speed-up of neutrino transformations in a supernova environment Sawyer, hep-ph/0503013
- 2. The multi-angle instability in dense neutrino systems Sawyer, arXiv:0803.4319
- 3. Neutrino cloud instabilities just above the neutrino sphere of a supernova Sawyer, arXiv:1509.03323
- 4. Self-induced neutrino flavor conversion without flavor mixing Chakraborty, Hansen, Izaguirre & Raffelt, arXiv:1602.00698
- 5. Fast pairwise conversion of supernova neutrinos: A dispersion-relation approach Izaguirre, Raffelt & Tamborra, arXiv:1610.01612
- Fast neutrino flavor conversions near the supernova core with realistic flavor-dependent angular distributions Mirizzi & Dasgupta, arXiv:1609.00528
- 7. Fast neutrino conversions: Ubiquitous in compact binary merger remnants Wu & Tamborra, arXiv:1701.06580

+ new developments ...

NEUTRINO ANGULAR DISTRIBUTIONS AT DECOUPLING

[Dasgupta, <u>A.M</u>., Sen, arXiV:1609.00528]



- Electron flavors remain in equilibrium with matter for a longer period than the non-electron flavors, due to the largest cross-sections of CC interactions
- Non-electron flavors decouple deeper in the star (more fwd-peaked distributions)
- Neutron-richness enhances CC interactions for v_e keeping them more coupled to matter (more isotropic distribution) than \overline{v}_{e} .

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FAST FLAVOR CONVERSIONS





FAST FLAVOR CONVERSIONS IN MULTI-D SNe

[Abbar et al., arXiV:1812.06883]



Which is the impact on the SN explosion?

Alessandro Mírízzí

NEUTRONIZATION BURST

[A.M., Tamborra, Janka, Saviano, Scholberg et al., arXiv:1508.00785 [astro-ph.HE]]



Robust feature of SN simulations

[Kachelriess et al., astro-ph/0412082, Gil-Botella & Rubbia, hep-ph0307244]



CONCLUSIONS

Observing SN neutrinos is the next frontier of low-energy neutrino astronomy

The physics potential of current and next-generation detectors in this context is enormous, both for particle physics and astrophysics.

Neutrino signal duration provides most useful particle-physics information. Neutrino signal from next nearby SN would make this argument much more precise.

Flavor conversions in SNe would provide valuable information on the neutrino mass hierarchy. Further investigations necessary on collective oscillations.

LOOKING FORWARD THE NEXT SN!



1-239[§]



La Rivista del Nuovo Cimento della Società Italiana di Fisica

Supernova neutrinos: production, oscillations and detection

A. MIRIZZI, I. TAMBORRA, H.-TH. JANKA, N. SAVIANO, K. SCHOLBERG, R. BOLLIG, L. HÜDEPOHL and S. CHAKRABORTY



arXiv:1508.00785 [astro-ph.HE]





<u>Home</u>

SN neutrinos at the crossroads: astrophysics, oscillations, and detection

From Monday, 13 May, 2019 - 08:00 to Friday, 17 May, 2019 - 14:00

Location: ECT* meeting room

Abstract:

The next Galactic supernova (SN) would be a once in a lifetime event for particle astrophysics, offering a unique opportunity for a multi-messenger detection of gravitational waves, neutrinos of all flavors, and multi-wavelength photons. The aim of the workshop is to review and discuss the recent advances in the study of SN neutrinos and future perspectives, ranging from SN simulations, neutrino oscillations, nucleosynthesis, gravitational waves, and SN neutrino experimental searches. The workshop is structured to address these broad issues in a synergistic interdisciplinary way and to stimulate discussions and collaborations among scientists working on SN neutrinos from different perspectives.

Registration period: 25 Feb 2019 to 22 Apr 2019

Website: https://indico.ectstar.eu/e/sn-neutrinos-at-the-crossroads_astrophysics-oscillat...

Organizers

Basudeb Dasgupta	Tata Institute for Fundamental Research	basudeb@gmail.com
Alessandro Mirizzi	University of Bari	alessandro.mirizzi@ba.infn.it

Attendees

No attendees yet.

Register Now 🗧

SECRETARIAT

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