# Phenomenological turbulent effects of core-collapse supernovae Toward to predict progenitor dependence

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# **1D** simulation Systematic study

Compactness:  $\xi_M = \frac{M/M_{\odot}}{R(M)/1000 \text{km}}$ 

$$M(r) = \int_0^r 4\pi \rho \, r'^2 \, dr'$$

↑ : Black hole (black)  $\xi_M$  $\downarrow$  : Supernova (Red)  $\xi_M$ 

**1D:** phenomenological simulation

### **Prediction :** property of progenitor like $\xi_M$ governs supernova explosion



<u> s19.8 (2002)</u>

s2014









# **3D** simulation



#### $M/M_{\odot}$ Compactness: $\xi_M$ *R*(*M*)/1000km

Progenitor $(M_{\odot})$	Envelope binding energy (10 <sup>51</sup> erg)	Compactness (calculated at 1.75 $M_{\odot}$ )			
s9.0	0.002	$3.831 \times 10^{-5}$			
s10.0	0.012	$2.165 \times 10^{-4}$			
s11.0	0.025	$7.669 \times 10^{-3}$			
s12.0	0.050	$2.215 \times 10^{-2}$			
s13.0	0.072	$5.932 \times 10^{-2}$			
s14.0	0.110	0.1243			
s15.0	0.144	0.1674			
s16.0	0.212	0.1546			
s17.0	0.251	0.1644			
s18.0	0.309	0.1715			
s19.0	0.341	0.1783			
s20.0	0.413	0.2615			
s25.0	0.865	0.3010			
s60.0	0.513	0.1753			

#### 1D prediction is Not consistent with latest results of 3D

### **One of the reasons : multi dimensional turbulent effects**



# Research of turbulence Turbulent convection in gain region



1.5 1.0 - 0.0 5.0 Specific Entropy [A.U.] 0.5 0.0 -1.0-1.5-2.0400

Turbulence depends on resolution, initial perturbations

Analyzing and understanding convection with multi- D simulation are difficult



### Analysis of Turbulence





non-linear analysis of turbulent effects is important to understand the mechanism.

Density perturbation triggers turbulence.

Kazeroni et al. 2018





# **Turbulent effects**



### Main turbulent effects (1) Turbulent Pressure $P_{\text{turb}} = \langle \rho' v' v' \rangle$

(2) Diffusion  $\langle e'v' \rangle$ 

(3) Dissipation  $\dot{e}_{\rm dis} = \rho v'^3 / L$ 



# Dimensionality (for Hydro)

3D

2D

1D

### Performance of selfconsistent simulation

Approximate transport

### Neutrino transport

Full Boltzmann

Our Goal Accuracy : 3D = 1D+ Resource : 3D >> 1D+

# Newtonian Full GR Gravity



# 1D+ simulation



1D+ is low computational cost. It is able to calculate neutrino heating and turbulent effects and to predict the progenitor dependence of CCSNe

Main turbulent effects are introduced self consistent 1D simulation

(1) Turbulent Pressure  $P_{\text{turb}} = \langle \rho' v' v' \rangle$ 

(2) Diffusion  $\langle e'v' \rangle$ 

(3) Dissipation  $\dot{e}_{\rm dis} = \rho v'^3 / L$ 

# Motivation

### Motivation : to predict progenitor dependence by using 1D+

### It needs to reproduce the 3D results with 1D+

### 1st step: researching turbulent effects in current 1D+



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# How to introduce turbulent effects Fluid equations: from 3D to1D

$$\frac{\partial \rho}{\partial t} + \nabla \cdot [\rho \overrightarrow{v}] = 0$$
$$\frac{\partial \rho \overrightarrow{v}}{\partial t} + \nabla \cdot [\rho \overrightarrow{v}^2 + P] = -\rho g$$
$$\frac{\partial (\rho e)}{\partial t} + \nabla \cdot [\overrightarrow{v} (\rho e + P)] = -\rho^2$$

#### Assuming spherical symmetry in the above three equations

mass conservation

### $g + S_{\nu}$ Euler equation

#### $\overrightarrow{v}g + Q_{\nu}$ Energy conservation



# How to introduce turbulent effects Reynolds decomposition (Murphy & Meakin 2011)

$$\frac{\partial \rho}{\partial t} + \nabla \cdot [\rho \overrightarrow{v}] = 0$$

$$\frac{\partial(\hat{\rho} + \rho')}{\partial t} + \nabla \cdot [(\hat{\rho} + \rho')(\hat{v} + \vec{v}')] = 0$$

$$\frac{\partial \langle \hat{\rho} + \rho' \rangle}{\partial t} + \frac{1}{r^2} \frac{\partial}{\partial r} r^2 \langle (\hat{\rho} + \rho')(\hat{v}_r + v_r') \rangle = 0$$

$$\frac{\partial \hat{\rho}}{\partial t} + \frac{1}{r^2} \frac{\partial}{\partial r} r^2 [\hat{\rho} \hat{v}_r + \langle \rho' v_r' \rangle] = 0$$

#### Phenomenological turbulent effect

Reynolds decomposition

 $A = \hat{A} + A'$ 

$$\hat{A} = \langle A \rangle = \frac{1}{4\pi} \int A(r, \theta, \varphi) \, d\Omega$$

$$\langle A' \rangle = 0$$



# Model: Mixing Length Theory

mix



# Scale of turbulence: $\Lambda_{mix}$ Turbulent velocity: v<sub>turb</sub>

Element is powered by buoyancy force









Mixing length parameters  

$$\Lambda_{\text{mix}} = \alpha_{\Lambda} H_{P} = \alpha_{\Lambda} \frac{P}{\rho g} \text{ mixing length}$$

Diffusion parameters

$$D_u = \alpha_u v_{\text{turb}} \Lambda_{\text{mix}} \quad (u = \epsilon, Y_e, K)$$



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# Explodability



$$\begin{array}{ll} \alpha_{\Lambda} & \text{Mixing length} & \text{parameters} \\ \Lambda_{\text{mix}} = \alpha_{\Lambda} H_P = \alpha_{\Lambda} \frac{P}{\rho g} & \text{mixing len} \end{array}$$

Progenitor : 9.6,12,15,20,25,30,35,40 *M*<sub>⊙</sub>

 $\alpha_{\Lambda} = 0.8 - 1.2$  total 8x17=136 models

 $12 M_{\odot}$  is difficult to explode.(Vartanyan et al.2018)

 $25 M_{\odot}$  is easy to explode.(Couch et al.2020)







### **Turbulent parameter dependence** Explosion energy



Turbulent parameter governs the progenitor dependence of  $E_{exp}$ 





# PNS mass (baryonic mass) **Turbulent parameter dependence**





Progenitor :12,15,20,25,30,35,40 *M*<sub>☉</sub>  $M_{\rm PNS}$  at  $t_{\rm b} = 0.5$  $M_{\rm PNS}\downarrow$  $\alpha_{\Lambda} \uparrow$  $\Rightarrow$ 



# compactness



ξ2.5

# Compactness : $\xi_M = \frac{M/M_{\odot}}{R/1000 \text{ km}}$

Progenitor :12,15,20,25,30,35,40*M*<sub>⊙</sub>

#### 40



## Dependence of compactness **Explosion energy**



Progenitor:12,15,20,25,30,35,40 *M*<sub>☉</sub>

$$E_{\rm exp}\left(t_{\rm b}=0.5\right)$$

Failure to explode :  $E_{exp} = 0$ ,

#### Turbulent parameter governs the compactness dependence of $E_{exp}$



# Dependence of compactness PNS mass (baryonic mass)



Progenitor :12,15,20,25,30,35,40 *M*<sub>☉</sub>

 $M_{\rm PNS}$  at  $t_{\rm b} = 0.5$ 



# Summary of progenitor dependence

- Explodability depends on turbulent parameter  $\alpha_{\Lambda}$ .
- Strong turbulence , which means large  $\alpha_{\Lambda}$  , trigger strong explosion.
- Focusing on  $\xi_M$ ,
- Compactness dependence of  $E_{\rm exp}$  is effected by turbulence parameter,  $\alpha_{\Lambda}$ .
- Compact dependence of  $M_{\rm PNS}$  is not effected by turbulence parameter,  $\alpha_{\Lambda}$ .



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## **Results** Evolution of Shock



# 1D+ has some tasks to reproduce 3D results accurately.

One is fitting the shock propagation.

In our 1D+, turbulent parameter effects prompt convection.

400



## Results **Turbulent velocity**



![](_page_27_Figure_2.jpeg)

200

![](_page_27_Picture_4.jpeg)

![](_page_27_Picture_5.jpeg)

### **Results** Source terms of turbulent energy in gain region

![](_page_28_Figure_1.jpeg)

$$\frac{\partial \rho v_{turb}^2}{\partial t} \dots = -\rho v_{turb}^2 \frac{1}{r^2} \frac{\partial (r^2 v_r)}{\partial r} + \rho v_{turb} \omega_{BV}^2 \Lambda_{mix} - \rho \frac{v_{turb}^3}{\Lambda_{mix}}$$

Source term of turbulent energy eq. :

$$S_{\text{turb}} = \rho v_{\text{turb}} \omega_{\text{BV}}^2 \Lambda_{\text{mix}}, \quad \epsilon_{\text{dis}} = |-\rho \frac{v_{\text{turb}}^3}{\Lambda_{\text{mix}}}$$

In our 1D+, dissipation term,  $\epsilon_{dis}$ , is large because of prompto convection.

![](_page_28_Picture_7.jpeg)

![](_page_28_Picture_8.jpeg)

# **Discussion of turbulent effects**

Phenomenological turbulent model of 1D+ is simple model. This is based on Mixing Length Theory(MLT).

However, turbulence of 3D depends on wave number. ( In our model, turbulence is evaluated by only one velocity,  $v_{\rm turb}\,$  )

Source of turbulent energy should be also corrected.

 $\chi = 20$   $\chi = 20$   $\chi = 20$   $\chi = 20$  0.6 0.4 0.2 4 0.2 4 0.2 4 0.2 4 0.5 10 10 10 15 20 25 $H k_x$ 

Kazeroni et al. 2018

![](_page_29_Picture_6.jpeg)

# **Next work** Analyzing turbulent effects with 3D

![](_page_30_Picture_1.jpeg)

![](_page_30_Picture_2.jpeg)

## Next work Analyzing turbulent effects with 3D

![](_page_31_Figure_1.jpeg)

![](_page_31_Figure_2.jpeg)

![](_page_31_Picture_3.jpeg)

# Summary

Motivation predicting progenitor dependence Today's talk : turbulence effects of 1D+ In our 1D+ model, Next steps are ...

- Results : turbulent parameter effects the progenitor dependence of  $\xi_M$

![](_page_32_Picture_5.jpeg)

Appendix

# Energy source of CCSNe

gravitational collapse. When a core of mass  $M_c$  collapses,  $E_g$  is

$$E_g = \left(-\frac{GM_c^2}{R_i}\right) - \left(\frac{GM_c^2}{R_f}\right) \sim \left(\frac{GM_c^2}{R_f}\right) = 3 \times 10^{53} \left(\frac{M_c}{M_\odot}\right)^2 \left(\frac{R_f}{10 \text{km}}\right)^{-1} \text{ [erg]}$$

 $R_{i,f}$ : initial (final) radius of core. ( $R_i > > R_f$ )

 $E_g$  is exchanged to internal energy, neutrino energy and kinematic energy. (Look 8 the slide of mechanism)

Energy source of a supernova explosion is the gravitational energy  $E_g$  released by

![](_page_34_Picture_8.jpeg)

![](_page_35_Figure_0.jpeg)

![](_page_35_Picture_2.jpeg)

![](_page_35_Picture_3.jpeg)

![](_page_35_Picture_4.jpeg)

### **Definition of** $E_{exp}, M_{Ni}, < E_{\nu} > , M_{PNS}$

$$E_{\rm exp} = \int_D \rho \epsilon + \rho v^2 / 2 + e_{\rm turb} + \rho \Phi \, \mathrm{d}V$$

$$M_{\rm Ni} = \int_D \hat{\rho} \hat{X}_{\rm Ni} \mathrm{d}V$$

$$\langle E_{\nu} \rangle_{tot} = \int_{t_b=0}^{t_b=0.5} L_{\nu} dt$$

 $M_{\rm PNS} = \int_{0}^{R_{\rm PNS}} 4\pi r^2 \rho dr \quad , R_{PNS} = R(\rho = 10^{11} \, [\rm g/cm^3])$ 

These properties are important to evaluate supernova mechanism.

# *dV* : Integral of non-binding region

#### $\hat{X}_{Ni}$ : mass fraction of nickel

#### $L_{\nu}$ : neutrino luminosity

![](_page_36_Picture_10.jpeg)

# **Appendix : calculation of** $f_{GW}$ with 1D+ Equation 3 in Sotani et al. 2021

1D+ can evaluate PNS mass  $M_{PNS}$  and PNS radius  $R_{PNS}$  from density distribution.

Finding the radius that is  $\rho = 10^{11}$  [g/c1

Derive the GW frequency  $f_{GW}$  from the following equation (Sotani et al. 2021)

 $f[kHz] = -1.410 - 0.443 \ln(x) + 9.337x - 6.7$ 

m<sup>3</sup>] 
$$\rightarrow R_{PNS} \quad M_{PNS} = \int_0^{R_{PNS}} 4\pi r^2 \rho dr$$

$$14x^2 , x \equiv \left(\frac{M_{\rm PNS}}{1.4M_{\odot}}\right)^{1/2} \left(\frac{R_{\rm PNS}}{10 \text{ km}}\right)^{-3/2}$$

## **Total emitted neutrino energy** Comparison with Compactness

![](_page_38_Figure_1.jpeg)

The progenitor dependence of  $\langle E_{\nu} \rangle_{tot}$  tends to be similar to compactness.

![](_page_38_Figure_3.jpeg)

![](_page_38_Picture_4.jpeg)

 $f^{t_b=0.5}$ 

 $\langle E_{\nu} \rangle_{tot} =$ 

![](_page_38_Picture_5.jpeg)

# Total emitted neutrino energy

![](_page_39_Figure_1.jpeg)

Progenitor :12,15,20,25,30,35,40 *M*<sub>☉</sub>

$$\langle E_{\nu} \rangle_{tot} = \int_{t_b=0}^{t_b=0.5} L_{\nu} dt$$

![](_page_39_Picture_4.jpeg)

## Frequency of GW **Comparison with Compactness**

![](_page_40_Figure_1.jpeg)

The progenitor dependence of  $f_{GW}$  tends to be similar to compactness.

![](_page_40_Figure_4.jpeg)

![](_page_40_Picture_5.jpeg)

## Frequency of GW Turbulent parameter dependence

![](_page_41_Figure_1.jpeg)

#### Progenitor:12,15,20,25,30,35,40*M*<sub>⊙</sub>

### $f_{\rm GW}\left(t_{\rm b}=0.5\right)$

 $\alpha_{\Lambda} \uparrow \quad \Rightarrow \quad f_{\rm GW} \downarrow$ 

![](_page_41_Picture_5.jpeg)

## **Correlation study** Table of correlation coeffici

The following is a test calculation of correlation coefficients

ient							
ZAMSmass –	1.000	0.869	0.901	0.881	0.847	0.259	0.876
ompactness –	0.869	1.000	0.985	0.987	0.950	0.397	0.972
PNSmass –	0.901	0.985	1.000	0.968	0.903	0.398	0.994
<e> -</e>	0.881	0.987	0.968	1.000	0.974	0.329	0.946
Eexp –	0.847	0.950	0.903	0.974	1.000	0.186	0.861
t(r_s=400) –	0.259	0.397	0.398	0.329	0.186	1.000	0.485
f_GW <b>-</b>	0.876	0.972	0.994	0.946	0.861	0.485	1.000
	ZAMSmass –	compactness –	PNSmass –	I A U V	Eexp –	t(r_s=400) –	f_GW -

![](_page_42_Figure_3.jpeg)

# **Correlation study Correlation with Total neutrino energy**

![](_page_43_Figure_1.jpeg)

### $M_{\rm PMS}$ has strongly positive correlation with $\langle E_{\nu} \rangle_{\rm tot}$ .

![](_page_43_Figure_4.jpeg)

# **Correlation study Correlation with Total neutrino energy**

![](_page_44_Figure_1.jpeg)

 $E_{exp}$  and  $M_{Ni}$  are positively correlated with  $\langle E_{\nu} \rangle_{tot}$ .

![](_page_44_Figure_4.jpeg)

But, these correlations depend turbulent parameter. 45

![](_page_44_Picture_6.jpeg)

# **Correlation study Correlation with GW frequency**

![](_page_45_Figure_1.jpeg)

![](_page_45_Figure_2.jpeg)

![](_page_45_Picture_3.jpeg)

# **Correlation study**

![](_page_46_Figure_1.jpeg)

But, these correlations depend turbulent parameter. 47

### $E_{exp}$ and $M_{Ni}$ are positively correlated with $f_{GW}$ .

![](_page_46_Picture_5.jpeg)

## Discussion Compared to 3D

![](_page_47_Figure_1.jpeg)

Our results of  $\alpha_{\Lambda} = 1.0$  is consistent to the result of 3D simulation

#### Results of 2D (Vartanyan et al. 2018)

![](_page_47_Figure_4.jpeg)

![](_page_47_Picture_5.jpeg)

![](_page_47_Picture_6.jpeg)

# **Discussion** Compared explodability to Couch et al. 2020

![](_page_48_Figure_1.jpeg)

ZAMS mass  $[M_{o}]$ 

Our feature is slightly different from previous work

However, the landscape of explodability is similar.

![](_page_48_Picture_5.jpeg)

![](_page_48_Picture_6.jpeg)

# Time dependence

![](_page_49_Figure_1.jpeg)

GW f mode [kHz]

#### the fitting formula of Gravitational wave (GW) by Sotani et al. 2021,

$$Hz) = -1.410 - 0.443 \ln(x) + 9.337x - 6.714x^2 ,$$

$$\left(\frac{M_{\rm PNS}}{1.4M_{\odot}}\right)^{1/2} \left(\frac{R_{\rm PNS}}{10 \,\rm km}\right)^{-3/2}$$

All properties of CCSNe depends on time.

Time dependence is also important factor for correlation study.

-> time-domain astronomy

![](_page_49_Picture_9.jpeg)

![](_page_49_Picture_10.jpeg)